## Note

# Accelerated Convergence of the Steepest Descent Method for Magnetohydrodynamic Equilibria

## 1. INTRODUCTION

Recently, a steepest-descent moment method [1] has been implemented for determining magnetohydrodynamic (MHD) equilibria. In Ref. [1] the Fourier amplitudes  $(R_{mn}, Z_{mn})$  of the cylindrical coordinates  $(R, \phi, Z)$  were discretized on a radial mesh. The resulting system of moment equations was solved using a descent algorithm to obtain the magnetic flux coordinate mapping.

In this paper, the  $\varepsilon$ -algorithm [2] is considered as a technique to accelerate the convergence rate of the steepest-descent method. The improved convergence rate of geometric sequences given by the  $\varepsilon$ -algorithm is well known. However, the relevance of this method specifically for the MHD equilibrium problem, and more generally in conjunction with the descent equations, has not been previously discussed.

The steepest-descent equations may be represented by the recursion relation:

$$O_n^{(i)} \mathbf{X}(n) = -\Delta W\{\mathbf{X}(n)\}.$$
 (1)

Here, X denotes the vector of discrete field amplitudes. For the MHD problem,  $X = \{R_{mn}, Z_{mn}\}$ . The  $\Delta$  notation symbolizes the gradient-difference operator for the positive definite energy functional W, and  $O_n^{(i)}$  is a first (i=1)- or second (i=2)-order linear opertor. Here, n assumes the role of a discrete time variable. The vector sequence X(n), n = 1,..., generated by the iteration of Eq.(1) may be decomposed as follows:

$$\mathbf{X}(n) = \mathbf{X}_{s}(n) + \mathbf{X}_{n'}(n) + \mathbf{X}_{\delta}(n).$$
<sup>(2)</sup>

Here, the dominant asymptotic limit of X(n) corresponds to the geometric vector sequence,

$$\mathbf{X}_{g}(n) = \mathbf{X}_{\infty} + \sum_{i=1}^{m_{x}} A_{j} r_{j}^{n} \mathbf{E}_{j}, \qquad (3a)$$

$$|r_j| < 1;$$
  $|r_1| \ge \cdots \ge |r_{m_x}|,$  (3b)

where  $\mathbf{X}_{\infty}$  is the stable equilibrium limit. The asymptotically vanishing nonlinear correction  $\mathbf{X}_{nl}$  is at least of second order,  $\mathbf{X}_{nl}(n \to \infty) \simeq O(|\mathbf{X}_g(n) - \mathbf{X}_{\infty}|^2)$ . The error term  $\mathbf{X}_{\delta}(n)$  contains the effects of round-off and truncation error. If this term

is random, it should not produce numerical stability problems in the application of the  $\varepsilon$ -algorithm.

In Eq. (3),  $r_j$  are the complex characteristic roots of the linearized version of Eq. (1). They are related through a dispersion relation to the eigenvalues  $\lambda_j$  of the Jacobian matrix  $\mathbf{J} \equiv -\nabla^2 W(\mathbf{X}_{\infty})$  evaluated at the equilibrium.  $E_j$  are the eigenvectors associated with  $\lambda_j$ . The eigenvalues of  $\mathbf{J}$  are negative for a stable equilibrium,  $\lambda_j \in [-|\lambda|_{\max}, -|\lambda|_{\min}]$ , and for the MHD energy functional W they satisfy [1] the stiffness property  $\delta_{\lambda} \equiv \lambda_{\max}/\lambda_{\min} \ge 1$ .

The goal of any acceleration scheme is to use a finite subset of the sequential data  $\mathbf{X}(n)$ , n = 1,..., generated by Eq. (1), to predict the desired equilibrium value  $\mathbf{X}_{\infty}$ . Consider the application of the scalar  $\varepsilon$ -algorithm to each component of the vector geometric sequence, Eq. (3a). A typical component of Eq. (3a) has the form  $(m \leq m_x)$ 

$$(\mathbf{X}_g)_s(n) \equiv S_n = S_\infty + \sum_{j=1}^m \alpha_j r_j^n.$$
(4)

It is known [3, 4] that for geometric sequences of the form given by Eq. (4) any successive 2m + 1 elements  $S_n, ..., S_{n+2m}$  can be used to determine  $S_{\infty}$  (the generalized Shank transform). However, because the theory makes use of Hankel determinants, it is impractical for computing with stiff systems. An alternative recursive nonlinear scheme [5], the  $\varepsilon$ -algorithm, also yields  $S_{\infty}$  after a finite number of operations involving 2m + 1 successive  $S_n$  members. Because of its recursive structure, the  $\varepsilon$ -algorithm is appropriate for stiff systems. The basic recursion relation is

$$\varepsilon_{k+1}^{(j)} = \varepsilon_{k-1}^{(j+1)} + \left[\varepsilon_k^{(j+1)} - \varepsilon_k^{(j)}\right]^{-1}, \qquad j = 0, \dots; \ k = 0, \dots,$$
(5a)

$$\varepsilon_{-1}^{(j)} = 0, \qquad \varepsilon_0^{(j)} = S_j, \qquad \text{for all } j. \tag{5b}$$

It can be proven [6] that if the parent geometric sequence satisfies certain minimal conditions then the k = 2m column,  $\varepsilon_{2m}^{(j)}$ , will have the limiting value  $S_{\infty}$ .

For stiff systems, a large number of iterations, N, of Eq. (1) is required to attain a single *e*-folding of the smallest characteristic root in Eq. (4), regardless of the order of  $O_n^{(i)}$ . In particular [1],  $N^{(i)} \simeq 0.5(\delta_{\lambda})^{1/i}$  for i = 1, 2. Alternatively, it is possible to use  $2m + 1 \leq N^{(i)}$  successive sequence members  $\{S_n\}$  in the  $\varepsilon$ -algorithm [Eq. (5)] to obtain  $S_{\infty}$  directly. Since many *e*-foldings are generally required to obtain  $S_{\infty}$  by direct iteration alone, the use of the epsilon algorithm in this context can be very economical.

Depending on the distribution of characteristic roots [for  $O^{(1)}$ ] or amplitudes [for  $O^{(2)}$ ] contributing to the parent sequence, considerable sequence acceleration may be manifested by columns of the  $\varepsilon$ -algorithm of even order less than 2m. This is apparent from the explicit form of the scalar algorithm of even order [7]

$$\varepsilon_{2k}^{(j)} = S_{\infty} + \frac{\sum_{\Omega_{k+1}^{m}} \prod_{l=1}^{k+1} (\alpha_{\rho_{l}} r_{\rho_{l}}^{j}) \prod_{1 \le p < q \le k+1} (r_{\rho_{p}} - r_{\rho_{q}})^{2}}{\sum_{\Omega_{k}^{m}} \prod_{l=1}^{k} [\alpha_{\rho_{l}} r_{\rho_{l}}^{j} (r_{\rho_{l}} - 1)^{2}] \prod_{1 \le p < q \le k} (r_{\rho_{p}} - r_{\rho_{q}})^{2}}.$$
 (6)

#### HANDY AND HIRSHMAN

Here,  $\Omega_k^m$  denotes the set of all distinct permutations of length k taken from a set of length m; and  $\rho_1$  is the subscript of the *l*th term in the particular permutation comprising  $\Omega_k^m$ . Note that  $\Omega_{m+1}^m$  is the null set, so that  $\varepsilon_{2m}^{(i)} = S_{\infty}$  as required.

For first-order systems, the fact that all of the characteristic roots in Eq. (3) have distinct moduli leads to the following asymptotic form for  $\varepsilon_{2k}^{(j)}$ :

$$\lim_{j \to \infty} \varepsilon_{2k}^{(j)} \sim S_{\infty} + O(r_{k+1}^j).$$
<sup>(7)</sup>

This estimate was obtained by noting that the dominant contributions to the numerator and denominator of Eq. (6) arise from terms containing the product of roots with the largest modulus. Recalling the root ordering in Eq. (3), it is apparent that each successive even  $\varepsilon$ -column converges faster than the preceding one and, in particular, faster than the parent sequence (the k=0 column of the  $\varepsilon$ -algorithm).

For second-order central difference schemes [1] the characteristic roots of Eq. (3) have the same modulus,  $r_j = \beta^{1/2} \exp(i\theta_j)$ , where  $\beta \sim 1 - \varepsilon$  and  $\varepsilon \sim \delta_{\lambda}^{-1/2} \ll 1$ . Although the theorem  $\varepsilon_{2m}^{(j)} = S_{\infty}$  still pertains, the asymptotic behavior of the even  $\varepsilon$ -columns will differ from that of first-order systems as given by Eq. (7). Nevertheless, some convergence improvement may be anticipated even for the  $2k < 2m \varepsilon$ -column, based on a smooth asymptotic approach (for increasing k) to the limiting value  $S_{\infty}$ . The numerical results in Section 3 also seem to confirm this expectation.

## 2. NUMERICAL STABILITY

Since the  $\varepsilon$ -algorithm is applied in practice to a sequence [Eq. (2)] that involves numerical perturbations of the exact asymptotic geometric sequence  $S_n$ , it is relevant to assess the numerical stability of the algorithm. One approach is to perturb an individual  $\varepsilon_{k+1}^{(j)}$  element and to generate the corresponding perturbations of  $\varepsilon_{k+1}^{(j)}$  and  $\varepsilon_{k+2}^{(j)}$  as given by the recursive algorithm, Eq. (5). In this manner [8] the following stability criterion for the relative error,  $\delta_k^{(j)} \equiv |\delta \varepsilon_k^{(j)}|/|\varepsilon_k^{(j)}|$ , is obtained (for  $S_{\infty} \neq 0$ ):

$$\lim_{j \to \infty} \delta_{2k+2}^{(j)} = R_{k+1} \delta_{2k}^{(j)},$$

$$R_{k+1} = |r_{k+1}|^2 / |1 - r_{k+1}|^2 \leq 1.$$
(8)

Clearly  $R_{k+1} \leq 1$  is desirable for stability and corresponds to either Re  $(r_k) \leq 0$  or  $|r_k| < \frac{1}{2}$ .<sup>1</sup> From Ref. [1], it is clear that  $|r_k|_{\max} \simeq 1$ , regardless of a first- or second-order formulation. Because of stiffness, it is also generally not possible to find explicit finite difference schemes for  $O_n^{(i)}$  in Eq. (1) with all characteristic roots satisfying  $-1 < \text{Re}(r_k) < 0$ . Nevertheless, by preconditioning the iterative scheme,

<sup>&</sup>lt;sup>1</sup> Also note that in addition to Eq. (8) one must also have the usual condition  $|r_k| < 1$ .

it may be possible [9] to obtain effective characteristic roots that do satisfy Re  $(r_k) < 0$  for all k.

If there were a preponderance of roots with negative real parts, then the majority of steepest-descent-generated sequences should be consistent with the sequence stability conditions. Some preliminary indications of this favorable root distribution are provided by the estimate  $|\lambda|_{\max}^{1/2} \sim \{\text{No. of radial mesh points}\}$  [1]. Thus, the eigenvalues of  $\nabla^2 W$  in Eq. (1) are expected to be packed closer to  $-|\lambda|_{\max}$  rather than to  $-|\lambda|_{\min}$ . The additional observation that for both first- and second-order systems the optimum  $O_n^{(i)}$  corresponds to roots satisfying<sup>2</sup> Re  $[r(-|\lambda|_{\max})] =$  $-\text{Re } [r(-|\lambda|_{\min})] < 0$  would support the expectation that the majority of steepest descent sequences are stable with respect to an  $\varepsilon$ -algorithm analysis.

An alternative interpretation of the above is that the  $\varepsilon$ -algorithm is more stable for nonmonotonic sequences, Re  $(r_k) < 0$ , than for monotonic ones. For the latter,  $R_{k+1} \leq 1$  can be satisfied by taking appropriate subsequences of the parent sequence. Consider  $S_n^* = S_{(In)}$ , n = 1, 2,... By selecting I so that

$$|r_j|_{\max}^{I} = \frac{1}{2},\tag{9}$$

or  $I \simeq 0.7 N^{(i)}$ , it is clear that the corresponding  $\varepsilon$ -algorithm will be numerically stable. The example presented in the next section corresponds to this situation.

## 3. NUMERICAL EXAMPLE

In Ref. [1] the following second-order operator was used in Eq. (1):

$$O^{(2)} \mathbf{X}_{(n)} \equiv \frac{\mathbf{X}(n+1) + \mathbf{X}(n-1) - 2\mathbf{X}(n)}{\Delta t^2} + \frac{\mathbf{X}(n+1) - \mathbf{X}(n-1)}{2\tau \Delta t}.$$
 (10)

The magnetic axis was treated accurately by adopting a Galerkin expansion for the m = 0 Fourier component of the radial inverse coordinate  $R(\rho, \theta, \phi)$ ,

$$R^{O_n}(\rho) = R_0^{O_n} + \sum_{k=1}^{\infty} u_k L_k(\rho^2), \tag{11}$$

where  $L_k$  are Legendre polynomials. Now consider the application of an  $\varepsilon$ algorithm analysis to the steepest-descent sequences for  $u_1$  and  $u_2$ . Two possible second-order theories have been considered here. One of these corresponds to keeping  $\tau$  and  $\Delta t$  fixed (nonoptimized situation). The other involves prescribing optimal  $\tau$  values [1]. It will be seen that application of the  $\varepsilon$ -algorithm to the first set of data, which proceeds to 3000 iterations, yields better results than the optimized  $\tau$ -

<sup>&</sup>lt;sup>2</sup> For first-order operators,  $r_k = 1 - |\lambda_k| \Delta t$ , where  $\Delta t$  is chosen optimally so that max ( $r_k$ ) is minimized with respect to k, i.e.,  $\Delta t^{\text{opt}} = 2/(|\lambda|_{\text{max}} + |\lambda|_{\text{min}}) \simeq 2/|\lambda|_{\text{max}}$ . Hence,  $r_k(-|\lambda|_{\text{max}}) \simeq -1$ . For second-order operators, a similar argument applies; see Handy [9].

varying situation at 4000 iterations. (Application of the  $\varepsilon$ -algorithm to the latter data set is invalid because of the variation of  $\tau$  during the iteration.) Table I corresponds to the nonoptimized ( $\tau = \text{const.}$ ) case. The parent sequence, up to 3000 iterations, is given in intervals of 100 iterations. At 3000 iterations,  $u_1 = -5.147 \times 10^{-2}$  and  $u_2 = -9.332 \times 10^{-3}$ . Because of the monotonic nature of the sequence, a subsequence analysis of the type discussed in the preceding section was implemented. Accordingly, using the expression  $I \simeq 0.7N^{(2)}$  and the estimate [1]  $N^{(2)} = 2\tau/\Delta t$  yields I = 90 for  $\Delta t = 0.04$  and  $\tau^{-1} = 0.4$ . The corresponding  $\varepsilon$ -algorithm analysis is represented in Table 3. Note the predicted values  $u_1 = -5.128 \times 10^{-2}$  and  $u_2 = -9.50e \times 10^{-3}$ .

These latter values for  $u_1$  and  $u_2$  compare well with those generated from an optimized ( $\tau$ -varying) second-order formulation as presented in Table II. At 3000 iterations,  $u_1^{opt} = -5.132 \times 10^{-2}$  and  $u_2^{opt} = -9.457 \times 10^{-3}$ , which are consistent with the  $\varepsilon$ -algorithm estimates. Indeed, the results contained in Tables II and III at 3000 iterations are far better converged than the corresponding entries of Table I. However, one can see that the results of the  $\varepsilon$ -algorithm actually surpass those of the optimized second-order code (Table II), because already on the basis of 3000 iterations of the Table I code the  $\varepsilon$ -algorithm predicts a reasonably stable limit value for both  $u_1$  and  $u_2$ . In contrast, the data in Table II show that  $u_2$  does not

Number of iterations	$u_1$	<i>u</i> <sub>2</sub>	
1000	$-5.475(\times 10^{-2})$	$-6.135(\times 10^{-3})$	
1100	-5.425	-6.716	
1200	-5.387	-7.163	
1300	- 5.353	-7.480	
1400	-5.323	- 7.748	
1500	- 5.294	- 7.997	
1600	-5.271	-8.208	
1700	-5.253	-8.386	
1800	-5.236	-8.542	
1900	-5.221	-8.677	
2000	- 5.208	- 8.791	
2100	- 5.197	-8.889	
2200	- 5.186	- 8.974	
2300	- 5.180	-9.047	
2400	- 5.173	-9.109	
2500	- 5.166	-9.162	
2600	-5.161	-9.208	
2700	- 5.157	-9.246	
2800	- 5.153	-9.279	
2900	5.150	-9.308	
3000	-5.147	-9.332	

TABLE I

## Nonoptimized, Second-Order Steepest-Descent Sequence for $u_1$ and $u_2$

begin to asymptote until at least 4000 iterations. There is some indication that the ultimate limit for  $u_2$  will be less than the  $-9.469 \times 10^{-3}$  entry and perhaps will even approach the  $\varepsilon$ -algorithm estimate of  $-9.50 \times 10^{-3}$ .

It will be noted from the data in Table III that the  $\varepsilon$ -algorithm data for  $u_1$  are slightly more susceptible to resonance effects than that of  $u_2$ . This is because the original parent sequence for  $u_1$ , as given in Table I, already converges much faster than that of  $u_2$ . Thus, considering Eq. (4) one can see that if two successive elements satisfy  $\varepsilon_k^{(j+1)} = \varepsilon_k^{(j)}$  then the ensuing recursively computed expression  $\varepsilon_{k+1}^{(j)}$  will be singular. Clearly, round-off error makes such infinities anomalously large. Theoretically, one can regulate away [10] such potential infinities. In practice, one

Number of				
iterations	<i>u</i> <sub>1</sub>	<i>u</i> <sub>2</sub>	1/τ	
1000	-5.488 (×10 <sup>-2</sup> )	$-6.006(\times 10^{-3})$	3.940 ( × 10 <sup>-1</sup>	
1100	-5.430	-6.624	3.060	
1200	- 5.368	-7.289	1.779	
1300	-5.306	-7.885	2.358	
1400	- 5.259	- 8.269	5.453	
1500	- 5.235	-8.477	1.351	
1600	-5.212	-8.728	2.084	
1700	- 5.191	-8.984	1.141	
1800	-5.161	-9.207	5.803	
1900	- 5.152	- 9.260	6.503	
2000	- 5.152	- 9.279	3.673	
2100	- 5.150	-9.315	3.970	
2200	- 5.144	-9.352	2.951	
2300	- 5.141	-9.376	3.165	
2400	- 5.139	- 9.395	3.665	
2500	- 5.137	-9.412	2.077	
2600	- 5.135	-9.428	1.518	
2700	- 5.134	- 9.441	4.211	
2800	-5.133	-9.450	4.007	
2900	-5.132	-9.454	1.312	
3000	-5.132	- 9.457	1.586	
3100	-5.132	-9.461	2.228	
3200	-5.132	9.464	5.863	
3300	- 5.131	-9.465	3.277	
3400	-5.131	9.466	5.838	
3500	- 5.131	-9.467	3.896	
3600	- 5.131	-9.468	3.311	
3700	- 5.131	-9.468	8.631	
3800	-5.131	- 9.469	6.143	
3900	-5.131	- 9.469	2.589	
4000	- 5.131	-9.469	2.798	

TABLE II

Optimized Second-Order, Steepest-Descent Parent Sequences for  $u_1$  and  $u_2$ 

may ignore them, provided they have no deleterious effects on successively higherorder columns of the  $\varepsilon$ -algorithm ansatz. This is clearly the behavior of both  $u_1$  and  $u_2$ , as given in Table III.

	$a$ -rigorithm as Applied to Data of fault 1, for $u_1$ and $u_2$ , respectively					
ε2	86	ε <sub>10</sub>	ε <sub>14</sub>	£ <sub>18</sub>		
$-5.243$ ( $\times 10^{-2}$ )						
- 5.064						
5.098	- 5.006					
-4.453	- 5.055					
- 5.183	- 5.139	-5.128				
-5.188	- 5.139	-5.128				
-4.947	- 5.112	- 5.128	-5128			
- 5.109	- 5.165	- 5.129	- 5.128			
- 5.124	$-3.002$ ( $\times 10^{5}$ )	- 5.130	- 5.128	- 5.128		
-5.137	3.603	- 5.122	- 5.128	-5.128		
- 5.148	-5.136	- 5.127	-5.128			
-5.116	-5.153	-5.128	-5.128			
-5.124	-5.087	- 5.129				
9.999 (×10 <sup>5</sup> )	-5.112	-5.127				
- 5.149	-5.134					
- 5.141	-5.135					
$7.206 (\times 10^6)$						
- 5.141						
$-8.654(\times 10^{-3})$						
-8.253						
-9.214	-9.190					
$-1.126(\times 10^{-2})$	-9.502					
-9.380	-9.523	-9.526				
-9346	-9.521	-9.509				
-9.648	-9.494	-9.485	- 9.491			
-9 545	-9.491	-9.489	-9.518			
-9410	-9487	-9.489	-9.485	- 9.499		
-9 500	-9.494	-9.489	-9.491	-9.510		
-9.530	-9479	-9.486	-9.492	2.210		
_ 0 <i>4</i> 01	-9485	-9484	-9490			
-9.458	- 9.489 - 9.488	-9481	7.770			
- 9.450 - 0.474	- 9.400 9.474	-9470				
-9.474	_9473	7.470				
0 107	- 9.470					
- 9.421	- 7.4/0					
- 9.49/						
- 7.310						

TABLE III

## ACKNOWLEDGMENTS

The authors thank J. C. Whitson and Professor R. Mickens for useful discussions. Research was sponsored by the Office of Fusion Energy, U.S. Department of Energy, under Contract W-7405-eng-26 and UN. CAR./OR 19X43371C(S/C) with the Union Carbide Corporation.

## References

- 1. S. P. HIRSHMAN AND J. C. WHITSON, Phys. Fluids 26 (1983), 3553.
- 2. A. GENZ, in "Pade Approximants" (P. R. Graves-Morris, Ed.), Institute of Physics, London, 1972.
- 3. D. SHANKS, J. Math. Phys. 34 (1955), 1.
- C. M. BENDER AND S. A. ORSZAG, "Advanced Mathematical Methods for Scientists and Engineers," McGraw-Hill, New York, 1978.
- 5. P. WYNN, Math. Tables Aids Comput. 10 (1956), 91.
- 6. J. B. McLEOD, Computing (Arch. Electron. Rechnen) 7 (1971), 17.
- 7. P. WYNN, SIAM J. Numer. Anal. 3 (1966), 91.
- 8. J. WIMP, "Sequence Transformations and Their Applications," Academic Press, New York, 1981.
- C. R. HANDY, "The Preconditioning of e-Algorithm Steepest Descent Methods for Linear Stiff Systems," Atlanta University preprint, 1984.
- 10. P. WYNN, Nord. Tidskr. Informat.-Behandl. 3 (1963), 175.

RECEIVED: April 17, 1984

C. R. HANDY

Department of Physics Atlanta University Atlanta, Georgia 30314

S. P. HIRSHMAN

Oak Ridge National Laboratory Oak Ridge, Tennessee 37830